Effects of local parity nonconservation in strong interactions in Pb-Pb collisions at LHC energy

V.N. Kovalenko

Saint-Petersburg State University

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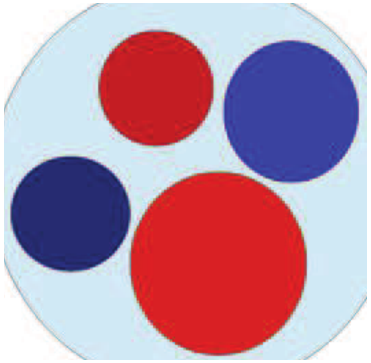
CP violation in QCD

$$\mathcal{L}_{\rm QCD} = -\frac{1}{4} G^{\mu\nu,a} G^a_{\mu\nu} + \bar{q}(i\gamma^\mu D_\mu - m)q,$$

$$D_\mu = \partial_\mu - igG^a_\mu \lambda^a, \quad G^a_{\mu\nu} = \partial_\mu G^a_\nu - \partial_\nu G^a_\mu + g f^{abc} G^b_\mu G^c_\nu$$

- θ-term $\qquad \Delta\mathcal{L}_\theta = \theta \frac{g^2}{16\pi^2} \mathrm{Tr}\left(G^{\mu\nu}\widetilde{G}_{\mu\nu}\right)$

- strong CP problem $\quad \theta \lesssim 10^{-9}.$
- P − and CP − odd bubbles may appear in a finite volume due to large topological fluctuations in a hot medium



- Gauge field configurations can be characterized by an integer topological (invariant) charge

A.A. Andrianov, V.A. Andrianov, D. Espriu, X. Planells, Phys. Lett. B 2012, 710, 230-235
A.A. Andrianov, V.A. Andrianov, D. Espriu, X. Planells, Proc. Sci., QFTHEP2013,025.
A.A. Andrianov, V.A. Andrianov, D. Espriu, X. Planells, AIP Conf.Proc. 1701 (2016) 1

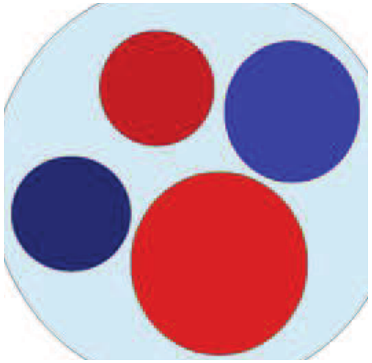
CP violation in QCD

- In QCD topologically non-trivial configurations of gauge fields can exist (instantons)

- Gauge field configurations can be characterized by an integer topological (invariant) charge

$$T_5 = \frac{g^2}{16\pi^2} \int_{t_i}^{t_f} dt \int_{\text{vol.}} d^3x \,\mathrm{Tr}\left(G^{\mu\nu} \widetilde{G}_{\mu\nu}\right) \in \mathbb{Z}$$

- Statistical treatment: with chemical potential μ_5



It must survive for a sizeable lifetime in a heavy-ion fireball,

$$\langle \Delta T_5 \rangle \neq 0 \quad \text{for} \quad \Delta t \simeq \tau_{\text{fireball}} \simeq 5 \div 10 \ \text{fm/c};$$

A.A. Andrianov, V.A. Andrianov, D. Espriu, X. Planells, Phys. Lett. B 2012, 710, 230-235
A.A. Andrianov, V.A. Andrianov, D. Espriu, X. Planells, Proc. Sci., QFTHEP2013,025.
A.A. Andrianov, V.A. Andrianov, D. Espriu, X. Planells, AIP Conf.Proc. 1701 (2016) 1

Chiral Magnetic Effect (CME)

Chiral Magnetic, Separation Effect:

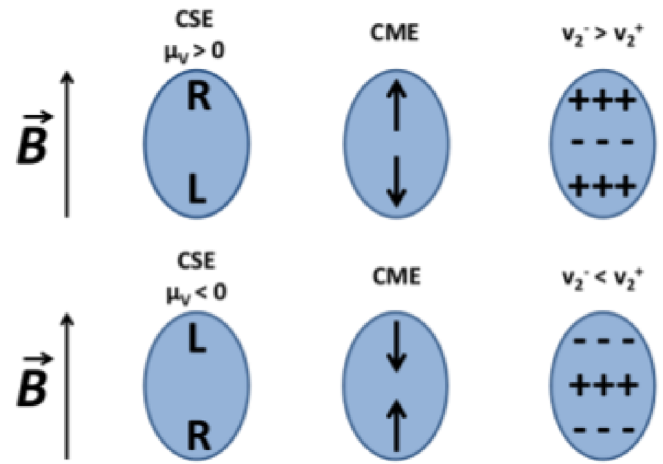
$$\vec{J}_V = \frac{N_c e}{2\pi^2}\mu_A \vec{B}, \quad \vec{J}_A = \frac{N_c e}{2\pi^2}\mu_V \vec{B}$$

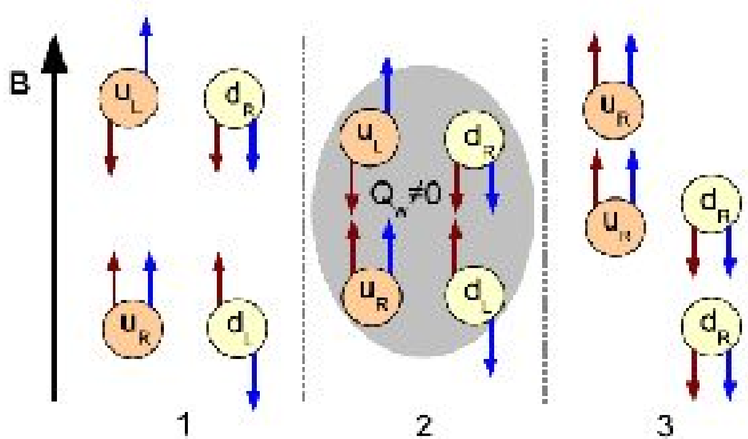
Thermodynamics:

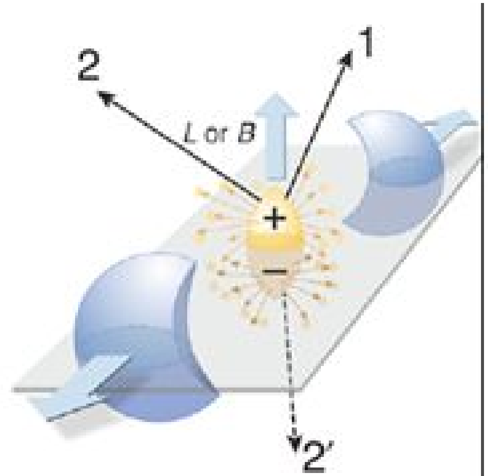
$$\vec{J}_V = \frac{N_c e}{2\pi^2}\chi\rho_A \vec{B}, \quad \vec{J}_A = \frac{N_c e}{2\pi^2}\chi\rho_V \vec{B}$$

Chiral basis:

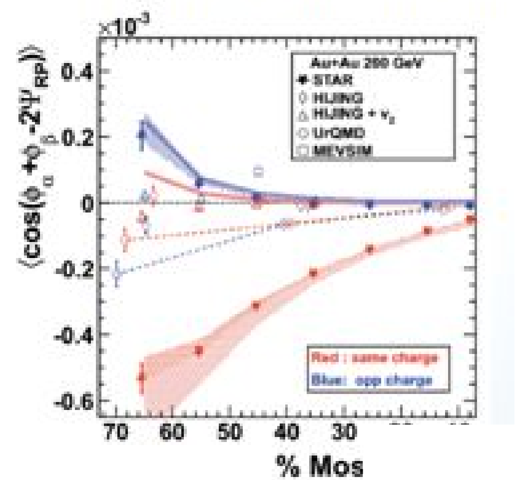
$$\vec{J}_L = -\frac{N_c e}{2\pi^2}\chi\rho_L \vec{B}, \quad \vec{J}_R = \frac{N_c e}{2\pi^2}\chi\rho_R \vec{B}$$







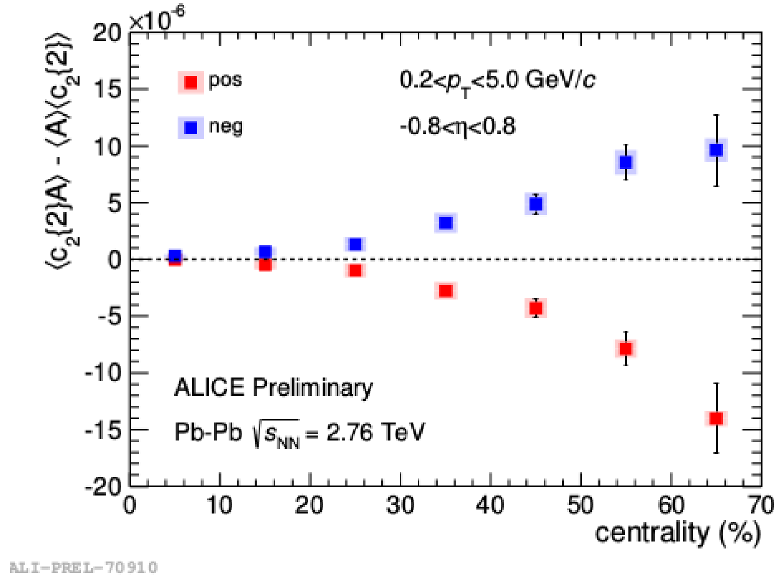
(STAR Collaboration)

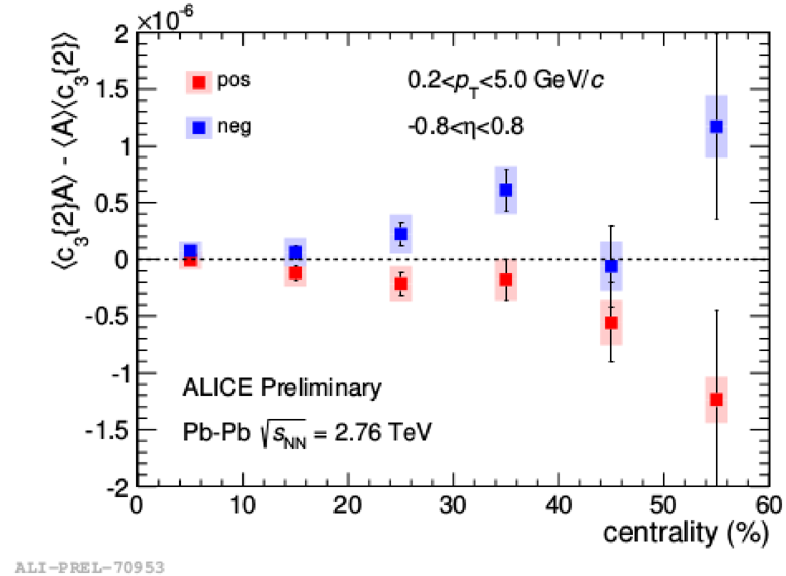


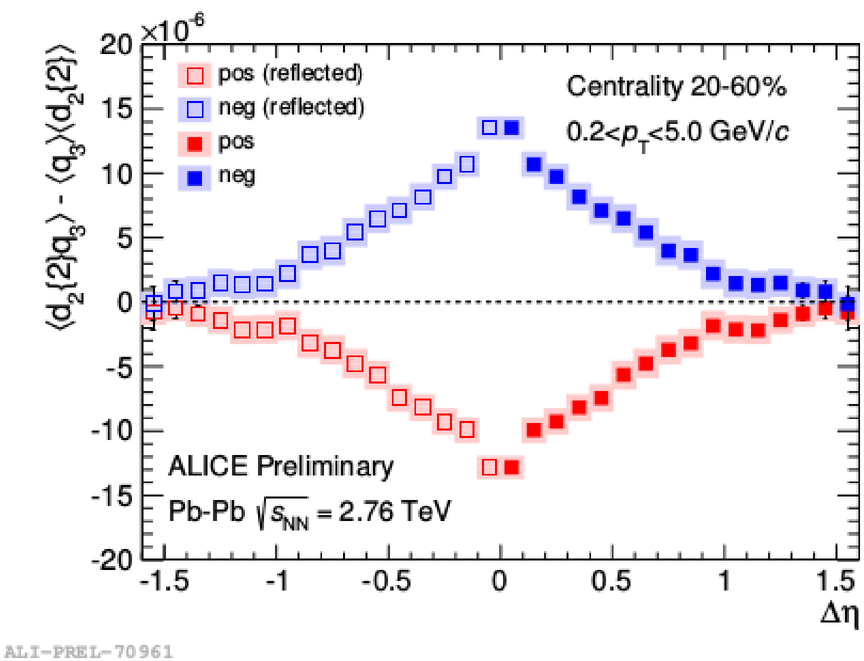
. Kharzeev, L.McLerran, H.Warringa, 2007

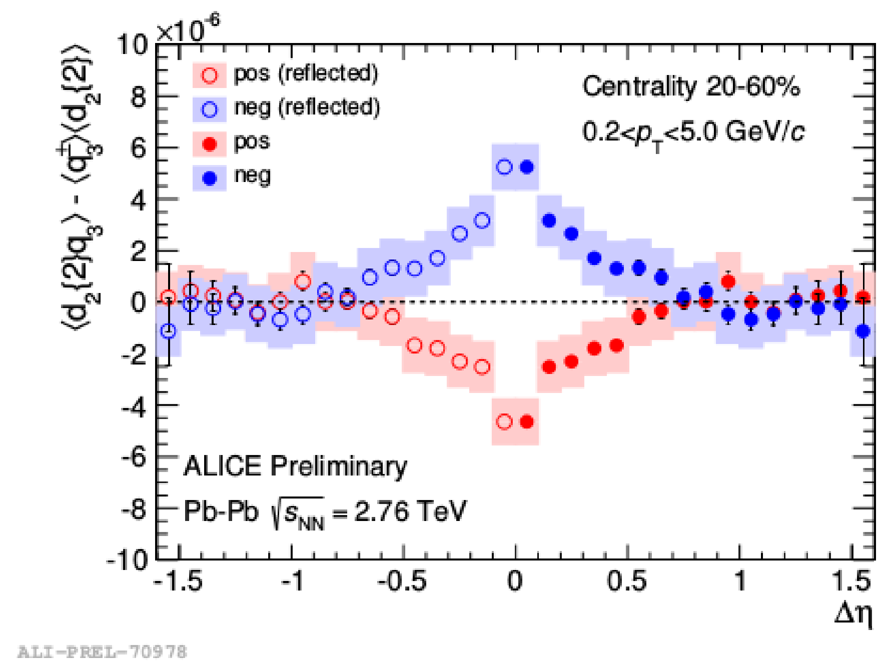
. Fukushima, D. Kharzeev, and H. Warringa. Phys. Rev. D, 78, 074033 (2008).

Chiral Magnetic Effect (CME) in ALICE









Jaroslav Adam , et al (ALICE Collab) Phys. Rev. C 93 (2016) 044903

New Possibilities

Parity forbidden decays

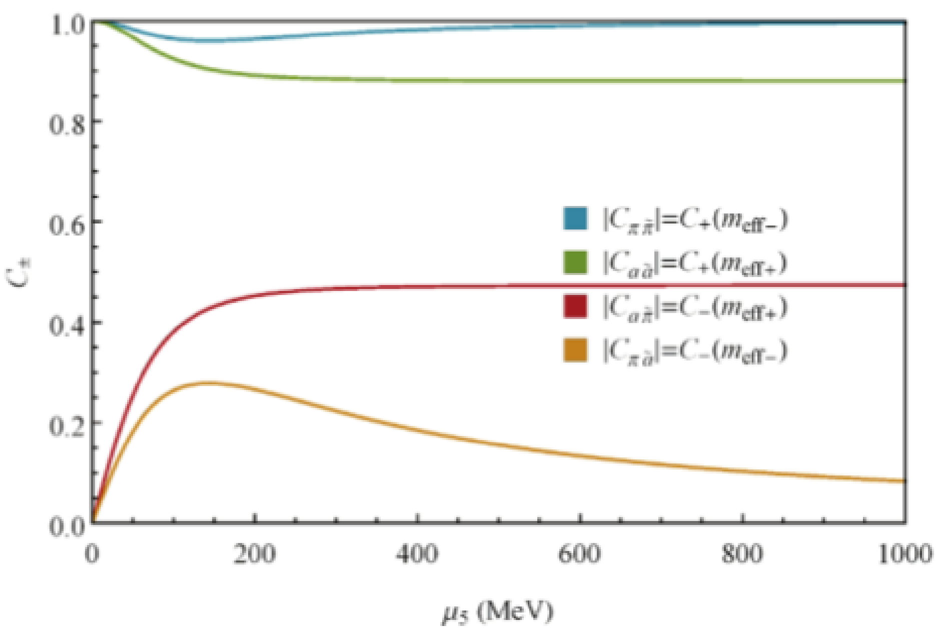
Effective meson theory in a medium with LPB

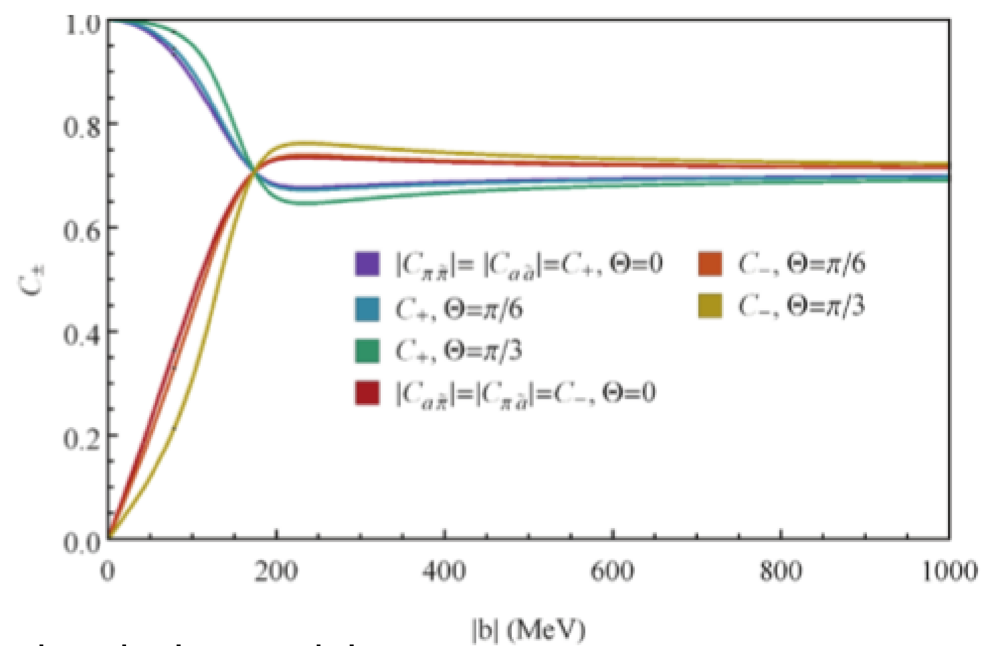
$$\mathcal{L} = \frac{1}{2}(\partial a_0)^2 + \frac{1}{2}(\partial \pi)^2 - \frac{1}{2}m_1^2 a_0^2 - \frac{1}{2}m_2^2\pi^2 - 4\mu_5 a_0 \dot{\pi},$$

$$m_1^2 = -2[M^2 - 2(3\lambda_1 + \lambda_2)v_q^2 - \lambda_2 v_s^2 - c v_s + 2\mu_5^2]$$

$$m_2^2 = \frac{2m}{v_q}B.$$

After diagonalization the new eigen-states appear: $\tilde{\pi}$ and $\tilde{a}_0$.





mixing coefficients dependence on chemical potential μ_5

A. A. Andrianov, V. A. Andrianov, D. Espriu, A.V. Iakubovich, A.E. Putilova EPJ Web of Conferences 158, 03012 (2017)

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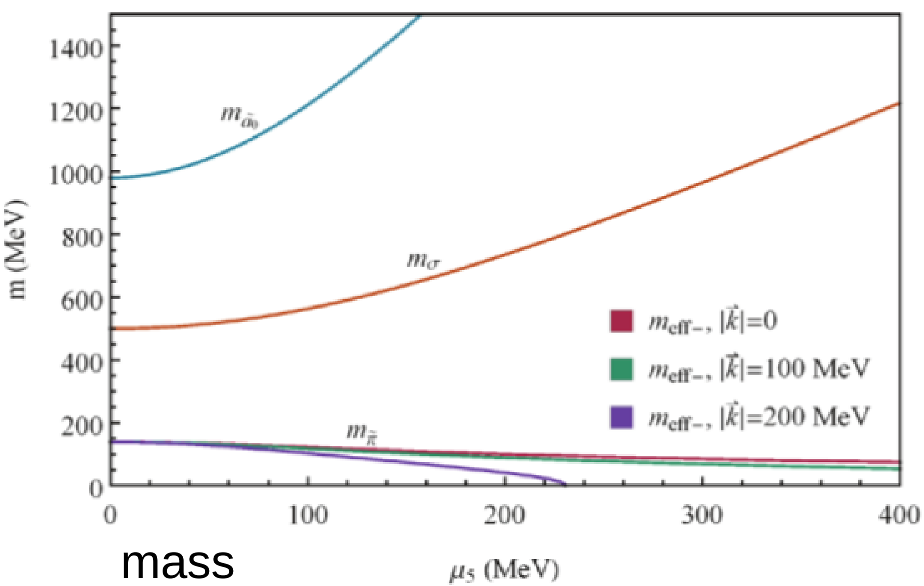
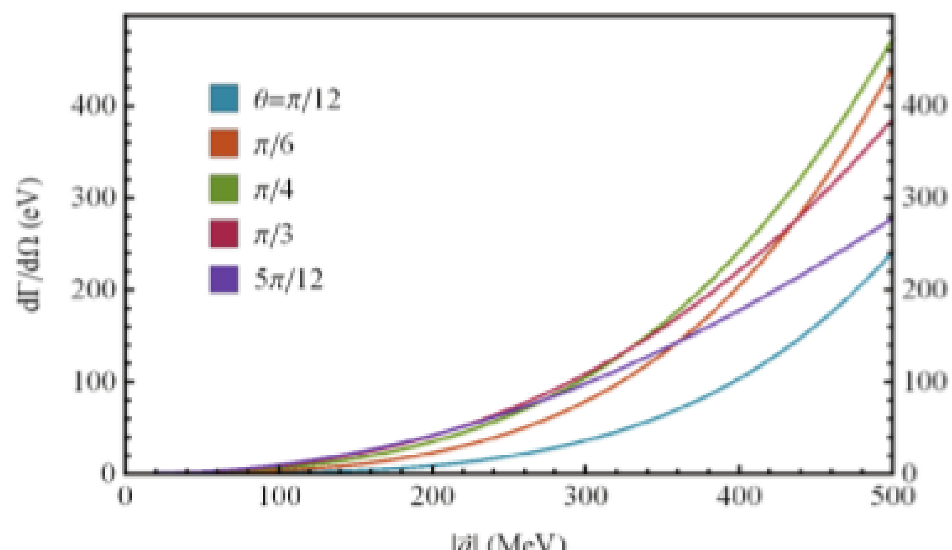
$$m_2^2 = \frac{2m}{v_q}B.$$

After diagonalization the new eigen-states appear: $\tilde{\pi}$ and $\tilde{a}_0$.

The decays $\tilde{a}_0^{\pm} \to \tilde{\pi}^{\pm}\gamma$.



mass



$\tilde{a}^{\pm}_0$ width

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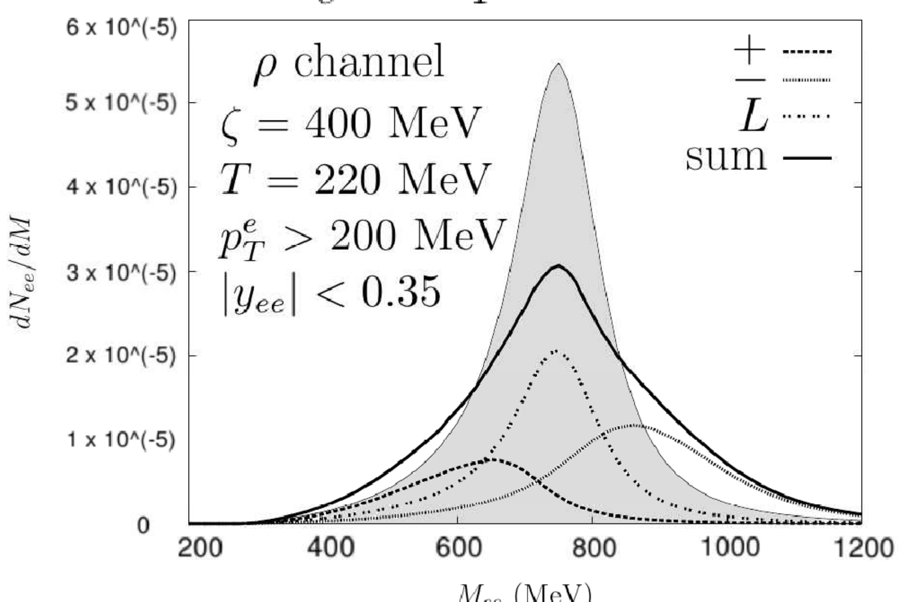
New Possibilities
DILEPTON POLARIZATION ANALYSIS
IN ρ ω → $l^+ l^-$ DECAYS

VDM + $\mathcal{L}_{\rm CS} = -\frac{1}{4}\varepsilon^{\mu\nu\rho\sigma}\,{\rm Tr}\left[\hat{\zeta}_\mu V_\nu V_{\rho\sigma}\right]$

The dilepton production from the $V(k) \to \ell^-(p)\ell^+(p')$ decays is governed by

$$\frac{dN_V}{dM} = \int \frac{d\tilde{M}}{\sqrt{2\pi}\Delta}\exp\left[-\frac{(M-\tilde{M})^2}{2\Delta^2}\right] c_V \frac{\alpha^2}{24\pi^2\tilde{M}}\Theta(\tilde{M}-n_V m_\pi)\left(1-\frac{n_V^2 m_\pi^2}{\tilde{M}^2}\right)^{3/2}$$
$$\times \int \frac{d^3\vec{k}}{E_k}\frac{d^3\vec{p}}{E_p}\frac{d^3\vec{p}'}{E_{p'}}\delta^4(p+p'-k)\sum_\epsilon \frac{m_{V,\epsilon}^4\left(1+\frac{\Gamma_V^2}{m_V^2}\right)}{\left(\tilde{M}^2-m_{V,\epsilon}^2\right)^2+m_{V,\epsilon}^4\frac{\Gamma_V^2}{m_V^2}}$$
$$\times P_\epsilon^{\mu\nu}(\tilde{M}^2 g_{\mu\nu}+4p_\mu p_\nu)\frac{1}{e^{\tilde{M}_T/T}-1}, \qquad \text{where } V=\rho,\omega \text{ and } n_V = 2, 0$$



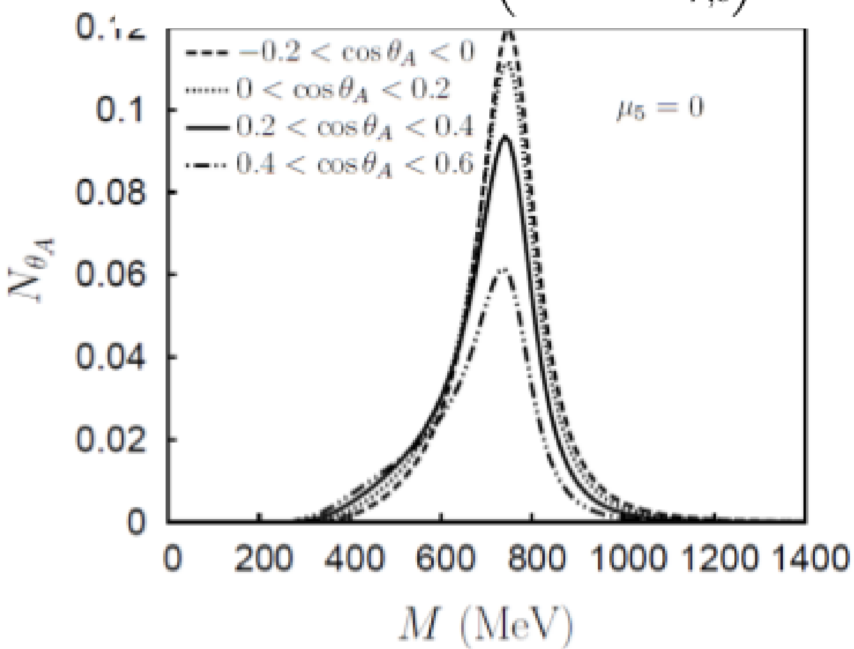
A. Andrianov, V. Andrianov, D. Espriu, EPJ Web of Conferences 137 , 01005 (2017)
Xumeu Planells, PhD thesis, November 2014, arXiv:1411.3283 [hep-ph]

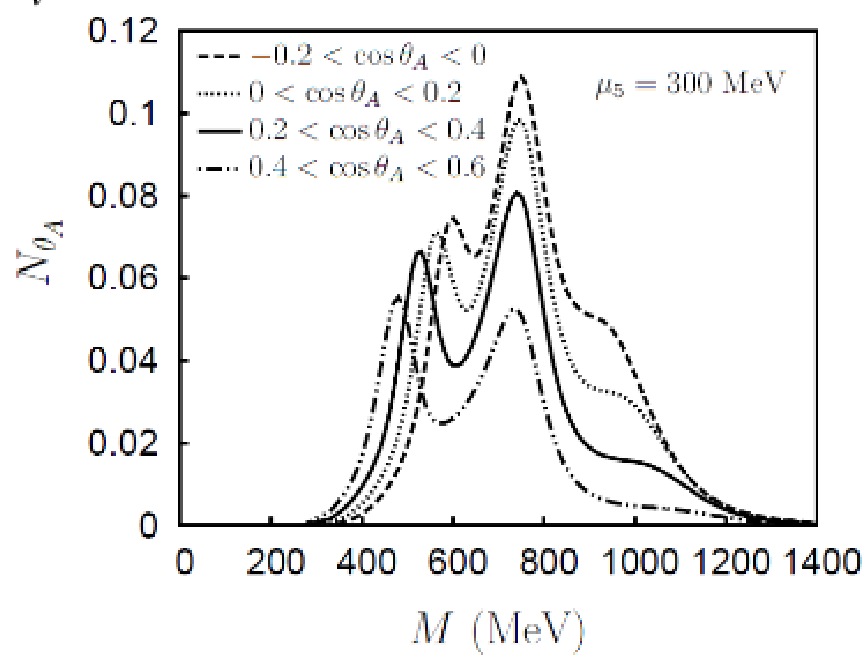
New Possibilities
DILEPTON POLARIZATION ANALYSIS
IN ρ ω → $l^+ l^-$ DECAYS

angle θ_A between the two outgoing leptons

$$\frac{dN_V}{dM d\cos\theta_A} = c_V \frac{\alpha^2}{6\pi M}\left(1 - \frac{n_V^2 m_\pi^2}{M^2}\right)^{3/2} \int \frac{p^2 p'^2 dp d\cos\theta d\phi}{E_p\sqrt{(M^2 - 2m_\ell^2)^2 - 4m_\ell^2(E_p^2 - p^2\cos^2\theta_A)}}$$

$$\sum_\epsilon \frac{m_{V,\epsilon}^4\left(1 + \frac{\Gamma_V^2}{m_V^2}\right)}{\left(M^2 - m_{V,\epsilon}^2\right)^2 + m_{V,\epsilon}^4\frac{\Gamma_V^2}{m_V^2}} P_\epsilon^{\mu\nu}(M^2 g_{\mu\nu} + 4p_\mu p_\nu)\frac{1}{e^{M_T/T} - 1},$$





A. Andrianov, V. Andrianov, D. Espriu, EPJ Web of Conferences 137 , 01005 (2017)
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New theoretical Results:

radiatively induced contribution to vector meson mass

$$\mathcal{L}_{\rm eff} = \mathcal{L}_{\rm kin} - \frac{1}{4}\varepsilon V_{\mu\nu}\varepsilon V^{\mu\nu} + \frac{1}{2}\varepsilon m^2 V_\nu V^\nu + \varepsilon\frac{b^\lambda b^\nu}{2b^2}V_{\lambda\rho}V_\nu^\rho - \frac{N_c}{8\pi^2}b_\mu V_\nu \epsilon^{\mu\nu\lambda\nu}V_{\lambda\nu}$$

Dispersion relations for ρ and ω mesons

$$\left\{k^2 - \tfrac{\xi}{m^2}\left[(b\cdot k)^2 - b^2k^2\right] - \bar{m}^2\right\}^2 - \zeta^2\left[(b\cdot k)^2 - b^2k^2\right] = 0$$

Mass of the transverse polarisations

$$m*^2 = \bar{m}^2 \pm \zeta b_0|\vec{k}| + \xi b_0^2|\vec{k}|^2/m^2$$

$$m_\omega^{*\,2} = (0.7766\mathrm{GeV})^2 + 0.9145\ b^2 \pm 2.7435\ b\ |\vec{k}| + 10.16(\mathrm{GeV}^{-2})\ b^2\ |\vec{k}|^2$$

$$m_\rho^{*\,2} = (0.7755\mathrm{GeV})^2 + 0.9119\ b^2 \pm 2.7357\ b\ |\vec{k}| + 10.13(\mathrm{GeV}^{-2})\ b^2\ |\vec{k}|^2$$

$$b_0 \ \simeq\ 0.5\mu_5,$$

Vladimir Kovalenko, Alexander Andrianov, Vladimir Andrianov,
J. Phys.: Conf. Ser. 1690, 012097 (2020), arXiv:2010.13238 [hep-ph]

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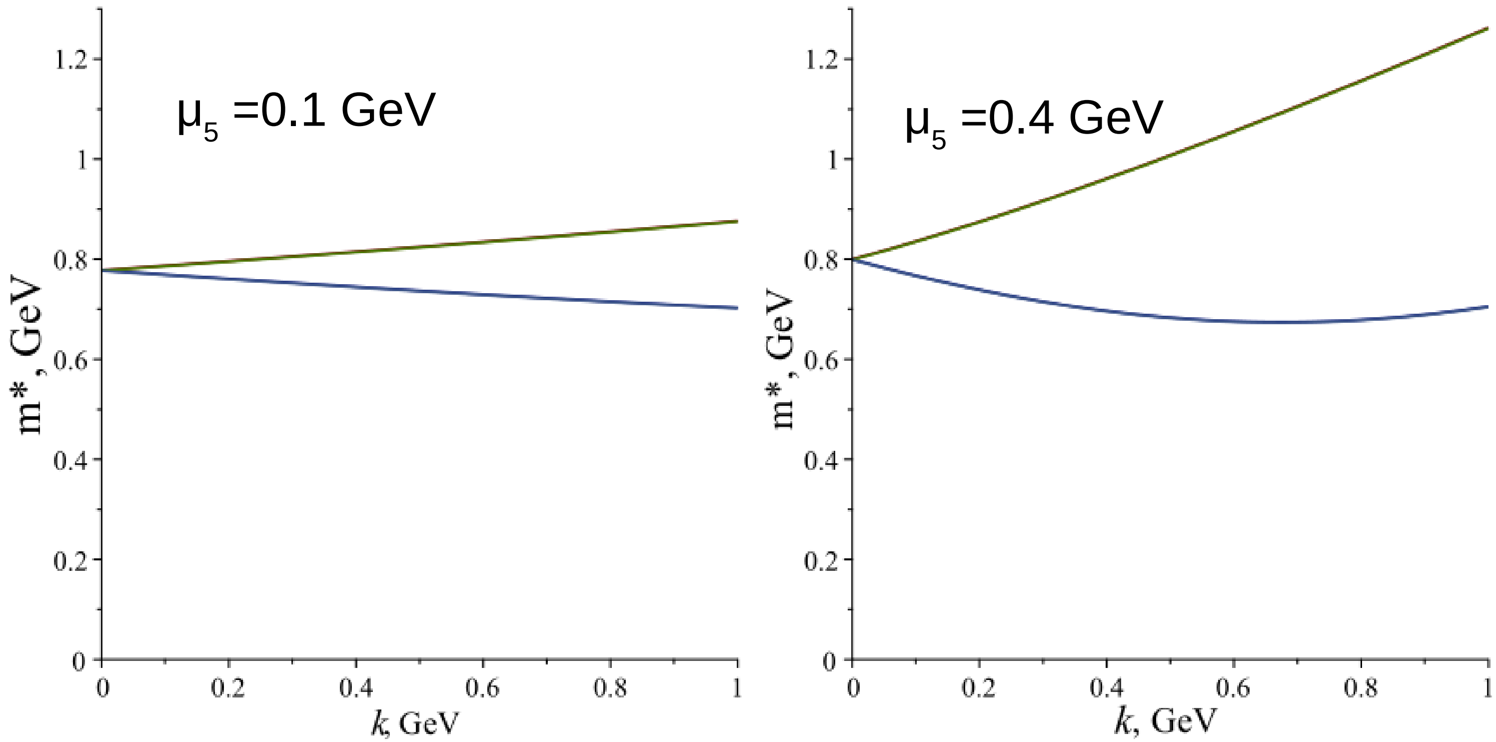
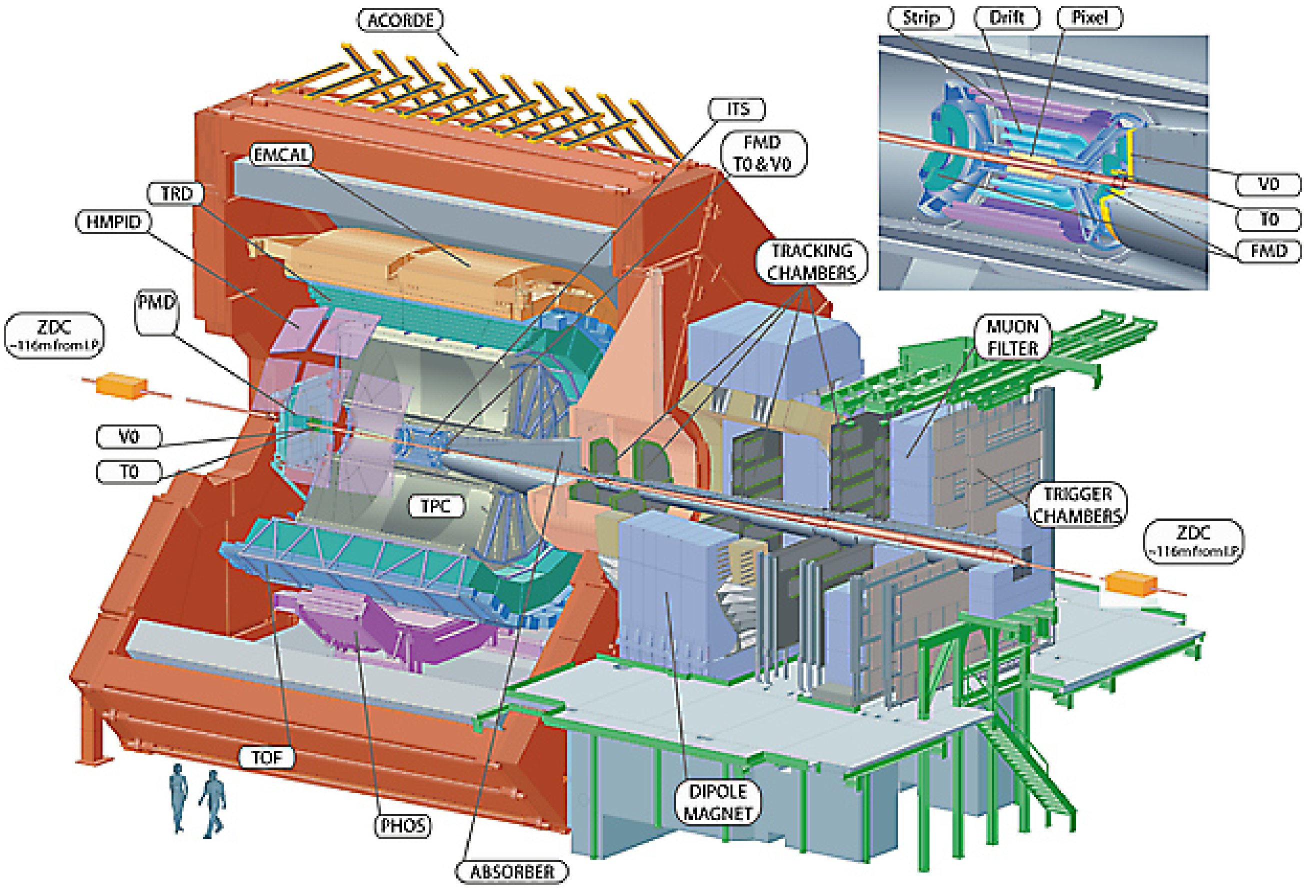


Vladimir Kovalenko, Alexander Andrianov, Vladimir Andrianov,
J. Phys.: Conf. Ser. 1690, 012097 (2020), arXiv:2010.13238 [hep-ph]

Experimental possibilities in heavy ion collisions at the LHC

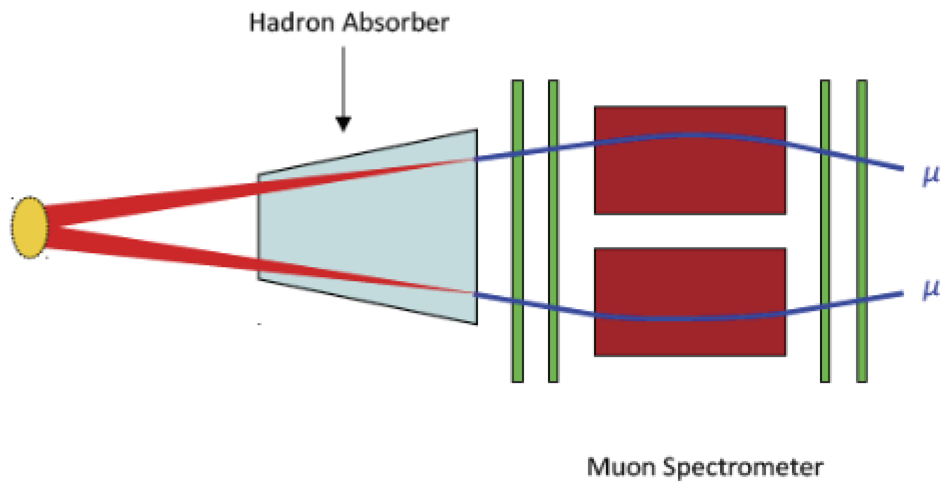


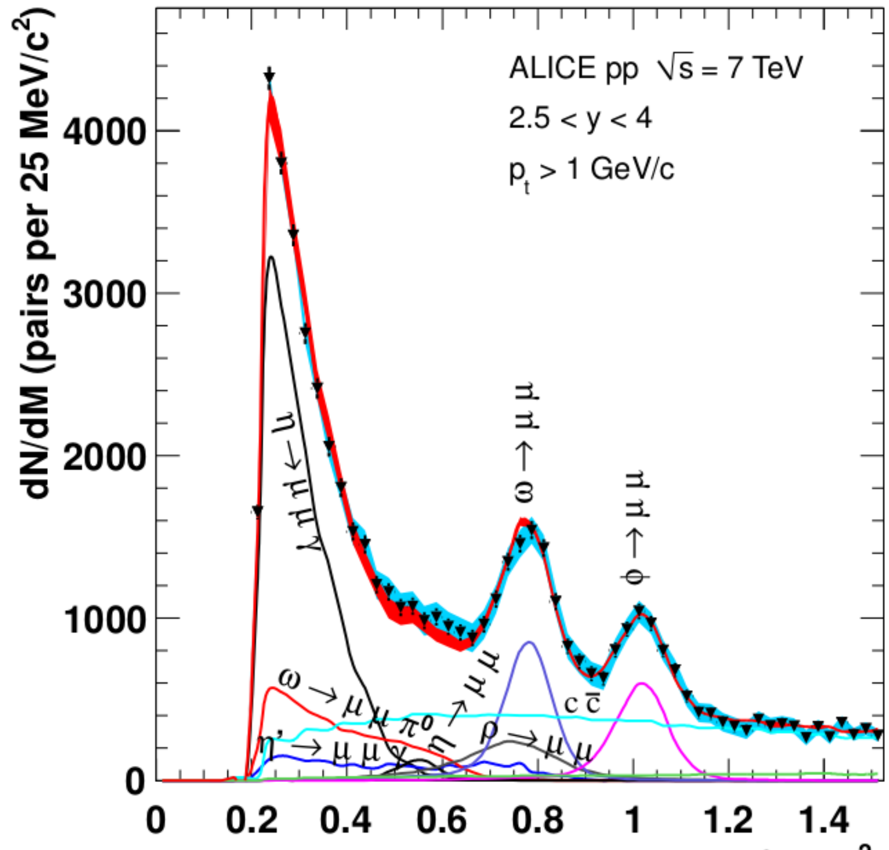
Detection systems

ITS+TPC, conversion method, EMCAL, PHOS, Muon Arm, PID

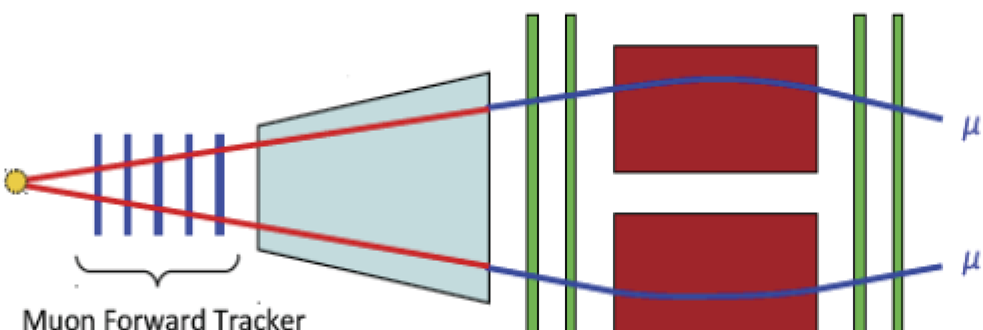
Di-muons

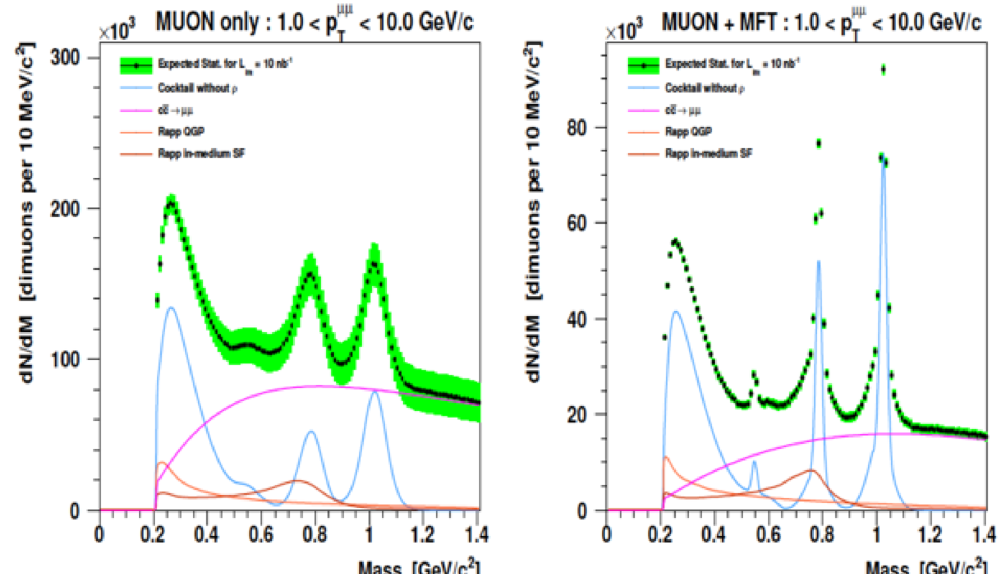
Muon Arm



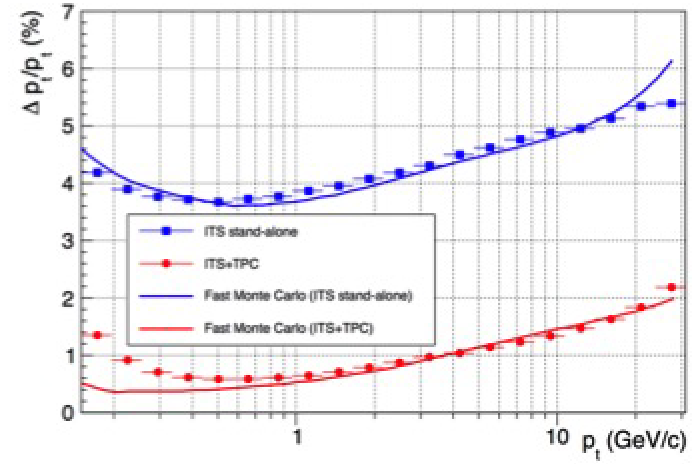


Phys. Lett. B710 (2012) 557-568





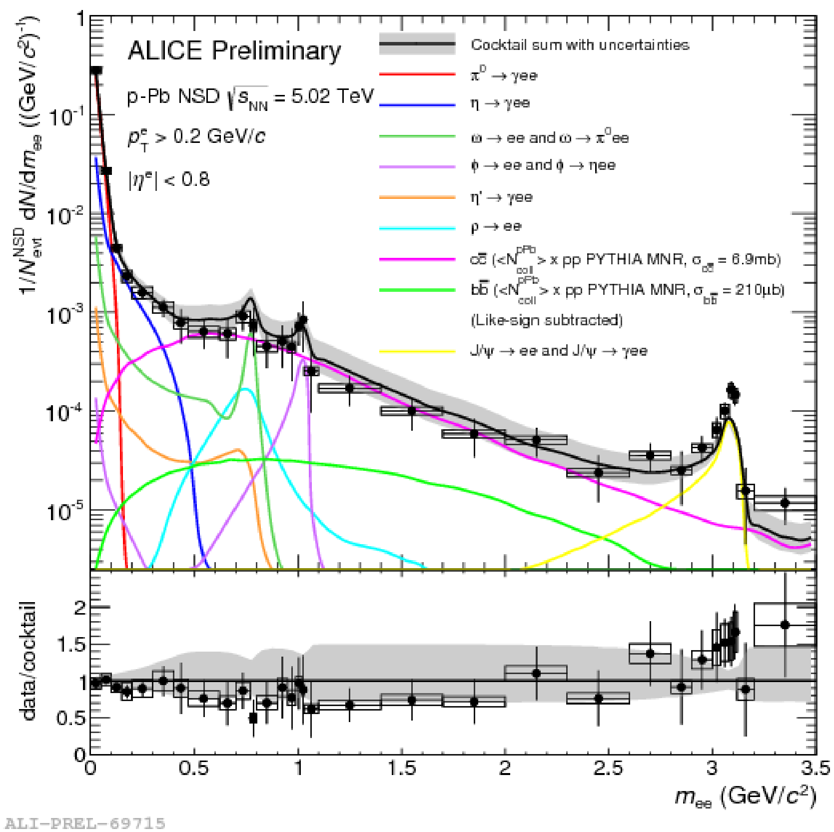
E. Garcia-Solis for the ALICE Collaboration, Nuclear and Particle Physics Proceedings 267–269 (2015) 382–391;
ALICE Collaboration, Technical Design Report for the Muon Forward Tracker, ALICE-TDR-018 (2015)



Di-electrons

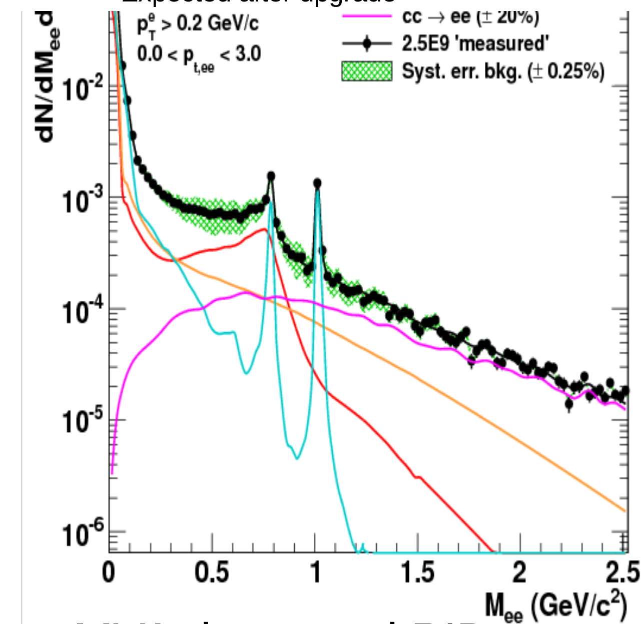
di-Lepton Detection

e+ e- in ITC+TPC, combied PID + MVA (ml) methods



Nucl.Phys. A956 (2016)

Expected after upgrade



MVA -improved PID

	Efficiency (%)	Purity (%)
Cut Method Efficiency Prioritised	70	91
Cut Method Purity Prioritised	13	99
Multivariate Method	95	96

https://indico.cern.ch/event/433345/contributions/2358113/attachments/1407816/2151898/poster.pdf

E. Garcia-Solis for the ALICE Collaboration, Nuclear and Particle Physics Proceedings 267–269 (2015) 382–391;
ALICE Collaboration, Technical Design Report for the Upgrade of the ALICE Inner Tracking System, J. Phys. G: Nucl. Part. Phys. 41 (2014) 087002

Monte Carlo setup

> Pythia 8.2 (Angantyr for heavy ion collisions, Pb+Pb, 5.02 TeV)

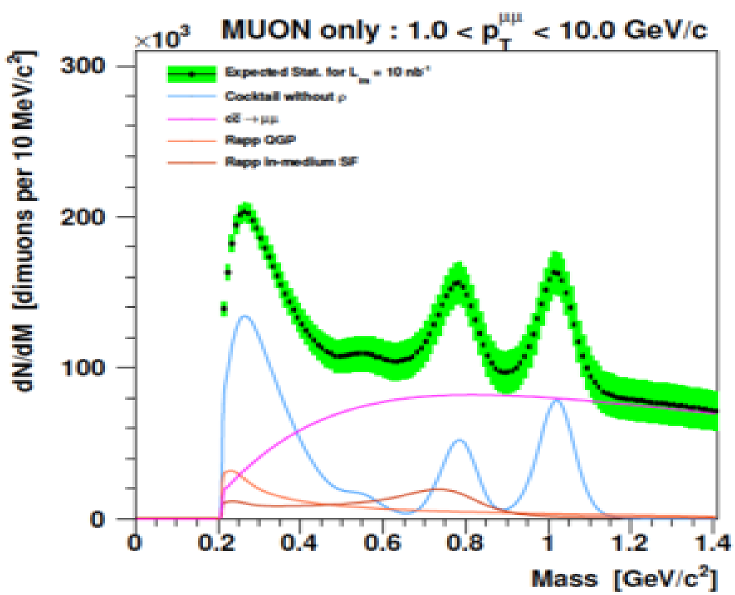
> Enhanced fraction of rho and omega leptonic decay channels

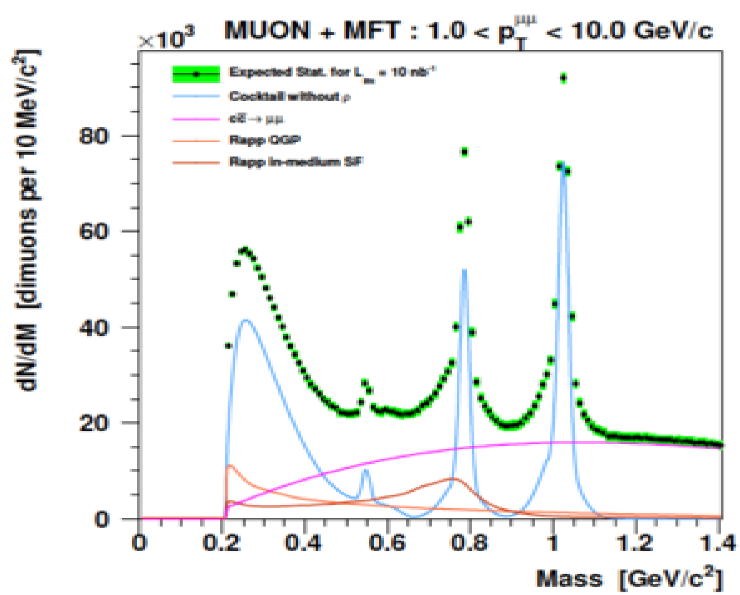
> Acceptance -0.8<eta<0.8 for di-Electrons, -3.6<eta<-2.45 for di-Muons

> Detector responce estimated using TDR resolutions/predictions (no fully detector modelling for this study yet)

> Focus on <u>resolution</u> of dimuon invariant mass studies (leaving significance/signal-over-background optimisation)

> Run 1+2 and Run 3 conditions

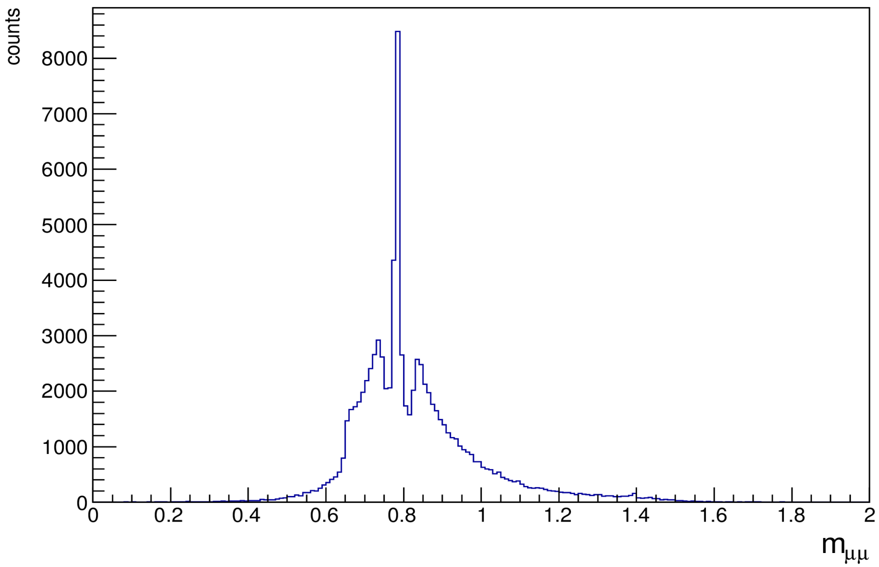


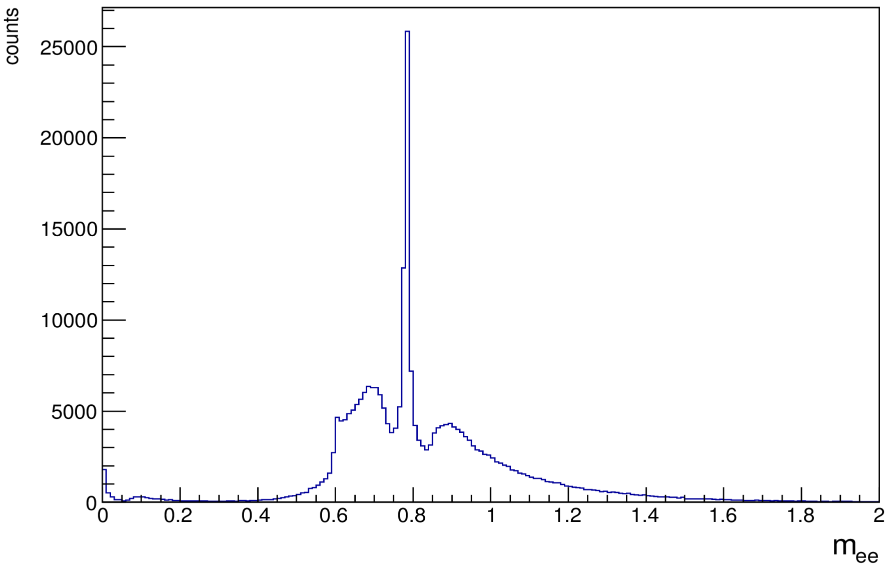


Pure Monte Carlo results (perfect detector response)

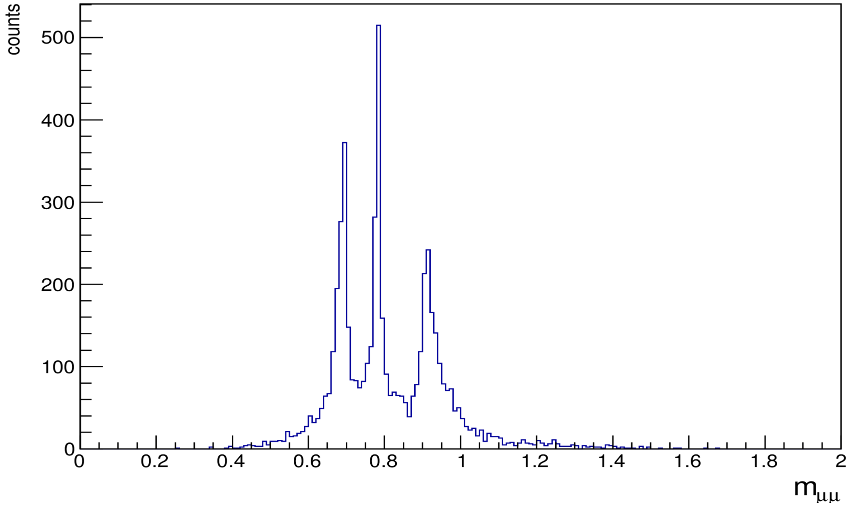
All: μ_5= 0.1 GeV

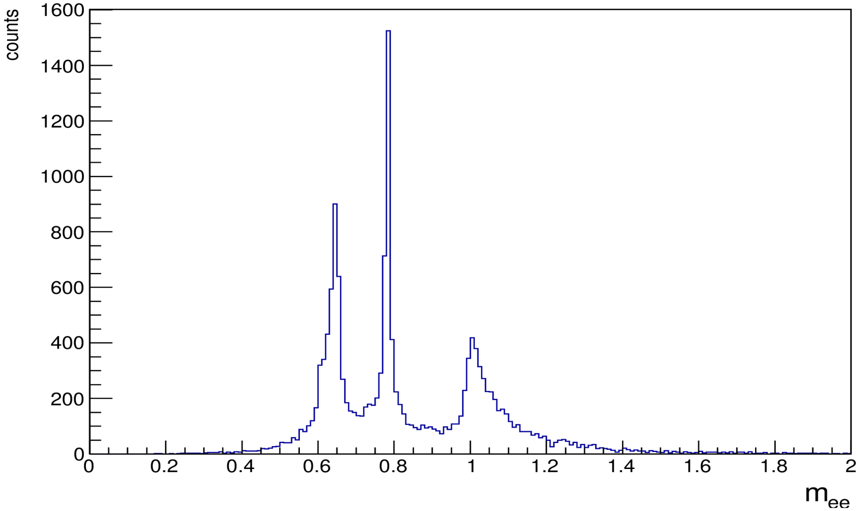
No angular θ_A selection





0.4<cos θ_A <0.5 (for muons — boost to midrapidity applied)





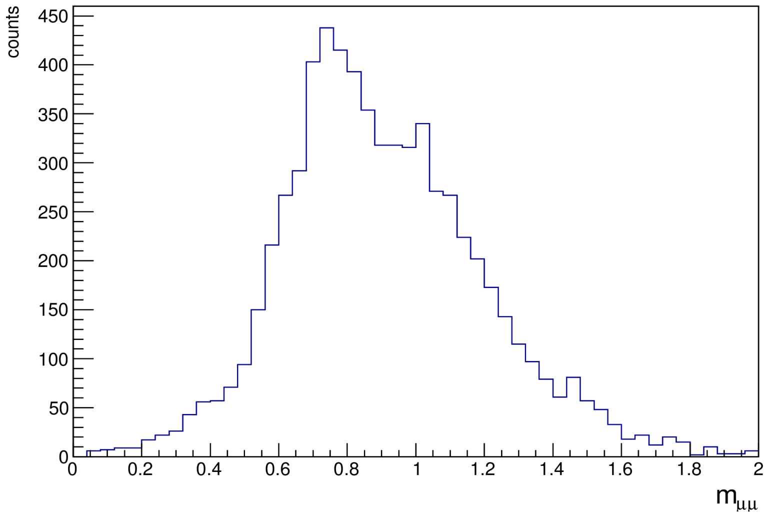
$p_{T\mu}$>0.2 GeV, $p_{T\mu\mu}$>0.4 GeV

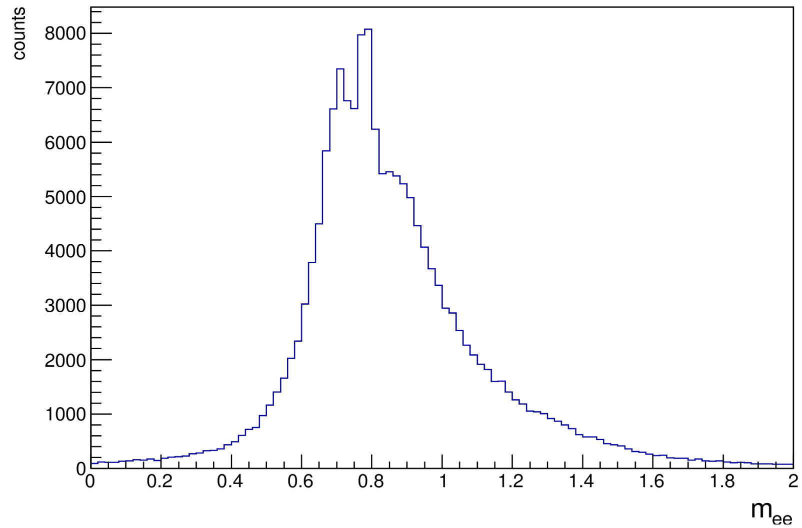
p_{Te}>0.1 GeV, p_{Tee}>0.4 GeV

Monte Carlo with smearing (Run 1+2 conditions)

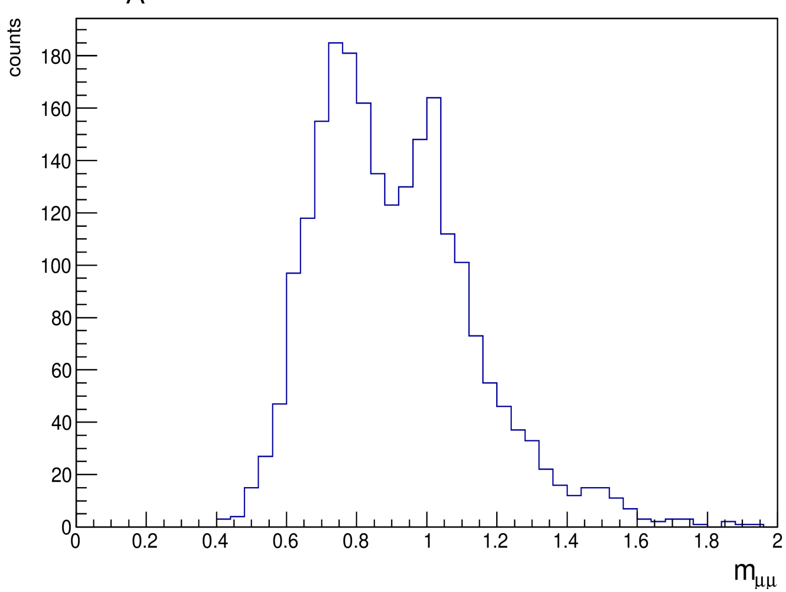
All: μ_5= 0.1 GeV

No angular θ_A selection



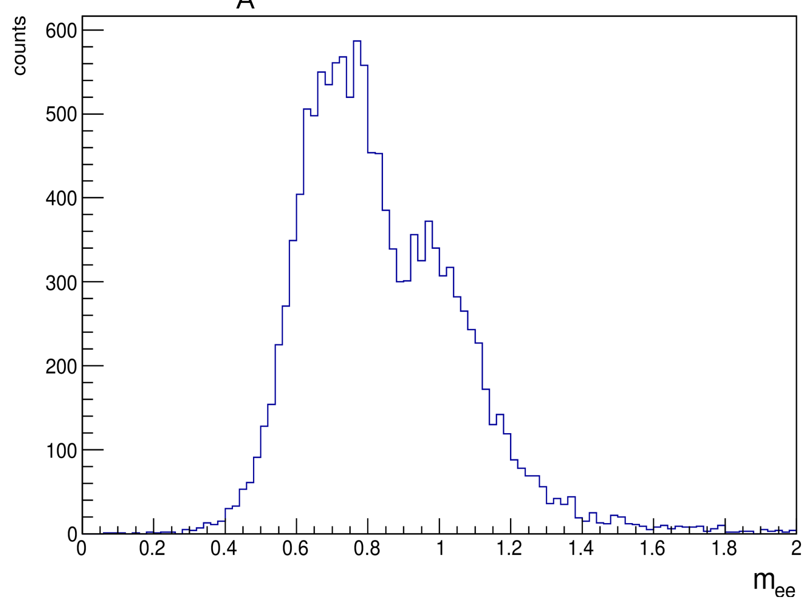


$0.4<\cos\theta_A<0.8$



$p_{T\mu}$>0.85 GeV, $p_{T\mu\mu}$>1.4 GeV

$0.4<\cos\theta_A<0.5$

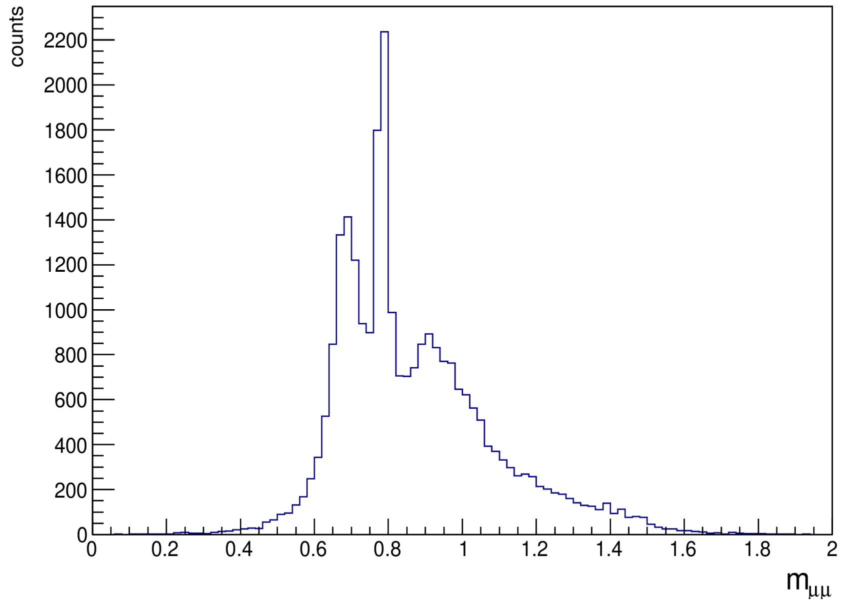


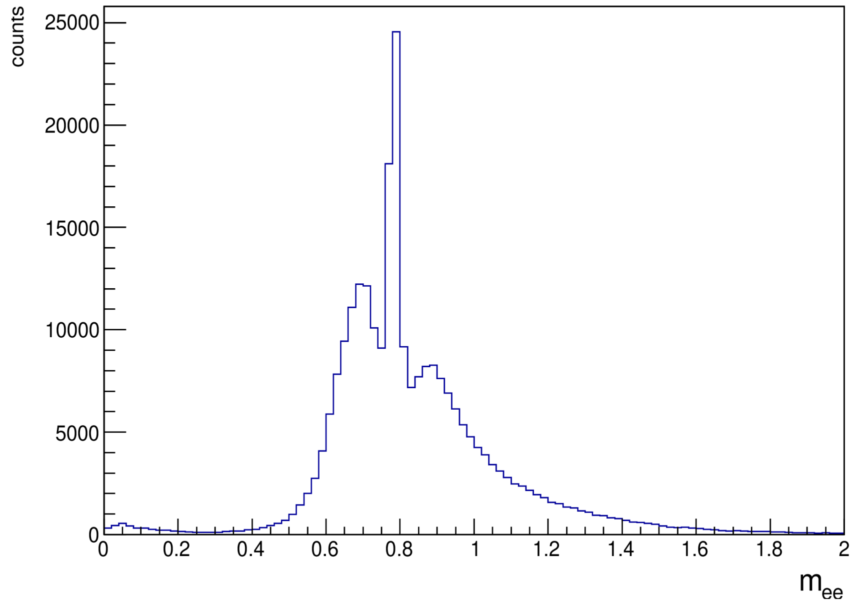
p_{Te}>0.3 GeV, p_{Tee}>0.4 GeV

Monte Carlo in Run 3 conditions

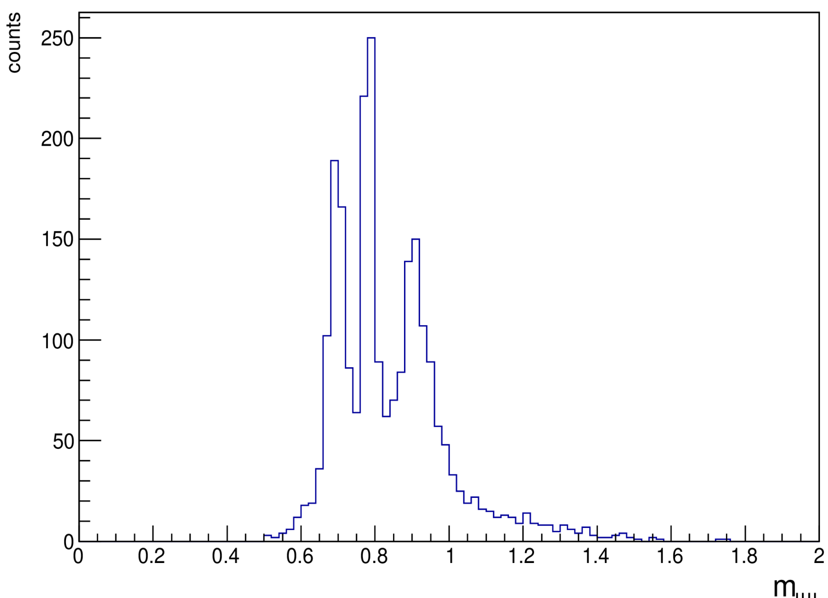
All: μ_5= 0.1 GeV

No angular θ_A selection



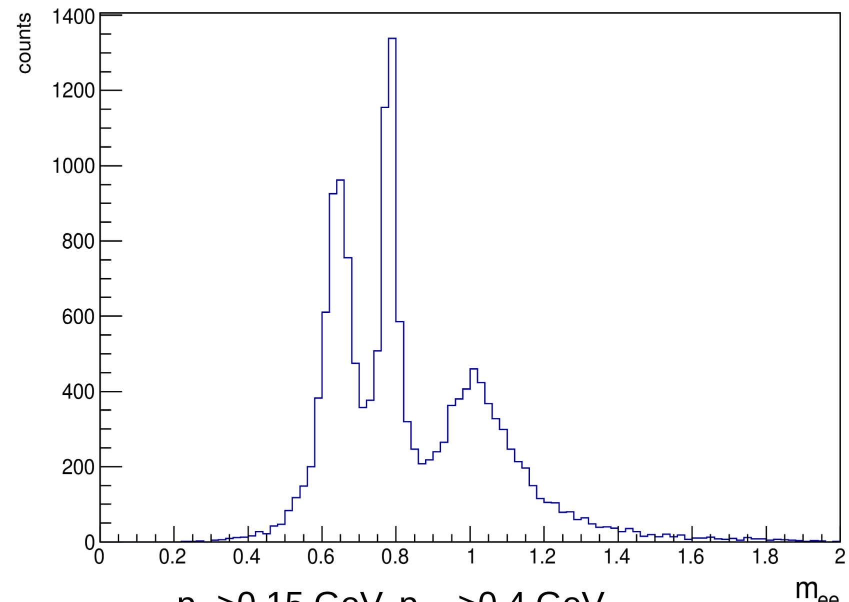


0.4<cos θ_A <0.5



$p_{T\mu}$>0.5 GeV, $p_{T\mu\mu}$>0.5 GeV

0.4<cos θ_A <0.5

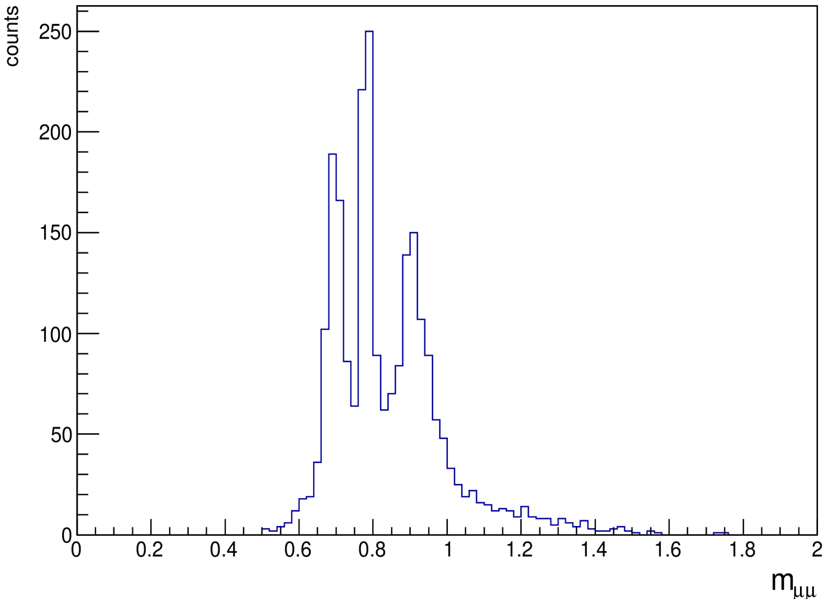


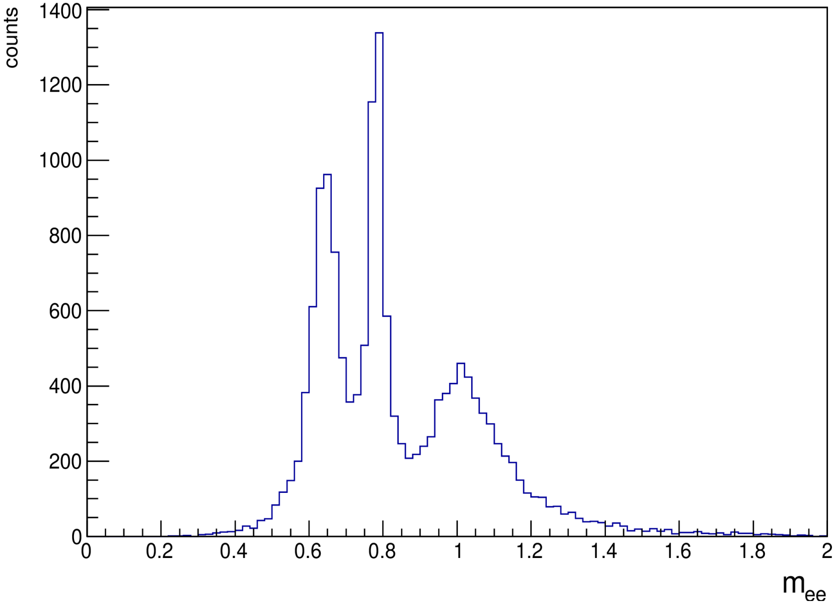
p_{Te}>0.15 GeV, p_{Tee}>0.4 GeV

Dependence on μ_5 (Run 3 conditions)

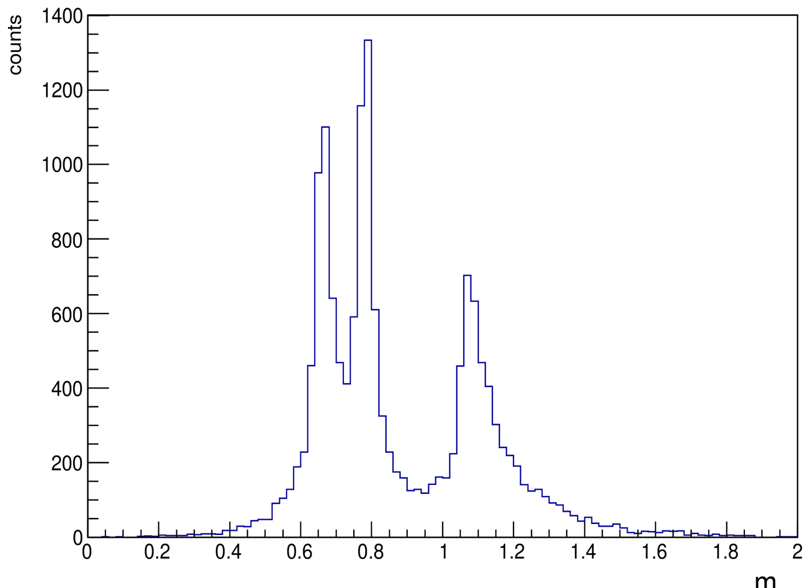
All: $0.4<\cos\theta_A<0.5$

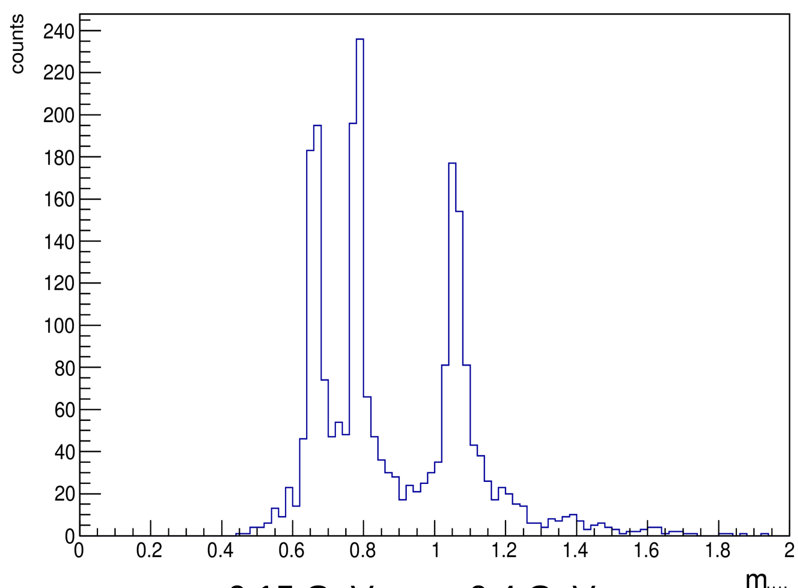
$\mu_5 = 0.1$ GeV





$\mu_5 = 0.2$ GeV





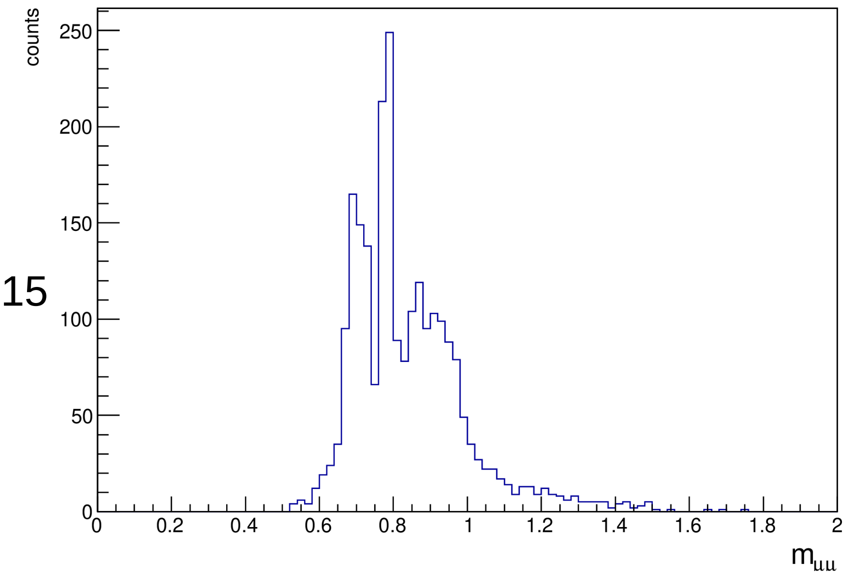
$p_{T\mu}>0.5$ GeV, $p_{T\mu\mu}>0.5$ GeV

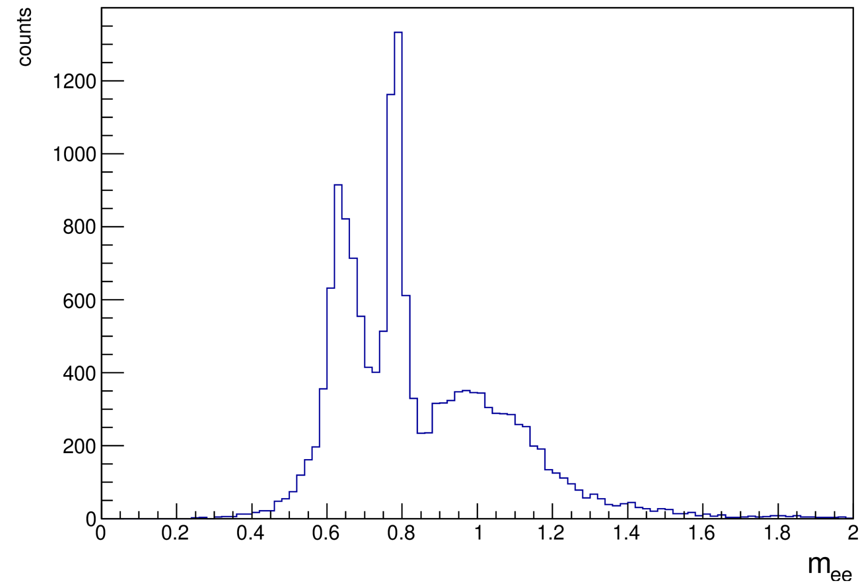
$p_{Te}>0.15$ GeV, $p_{Tee}>0.4$ GeV

Influence of the fluctuation of μ_5 (Run 3 conditions)

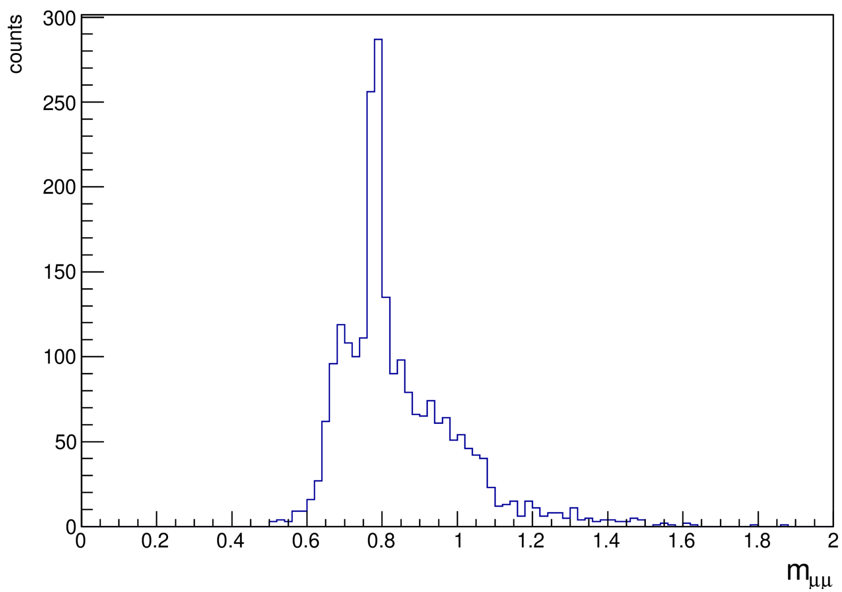
All: $0.4<\cos\theta_A<0.5$

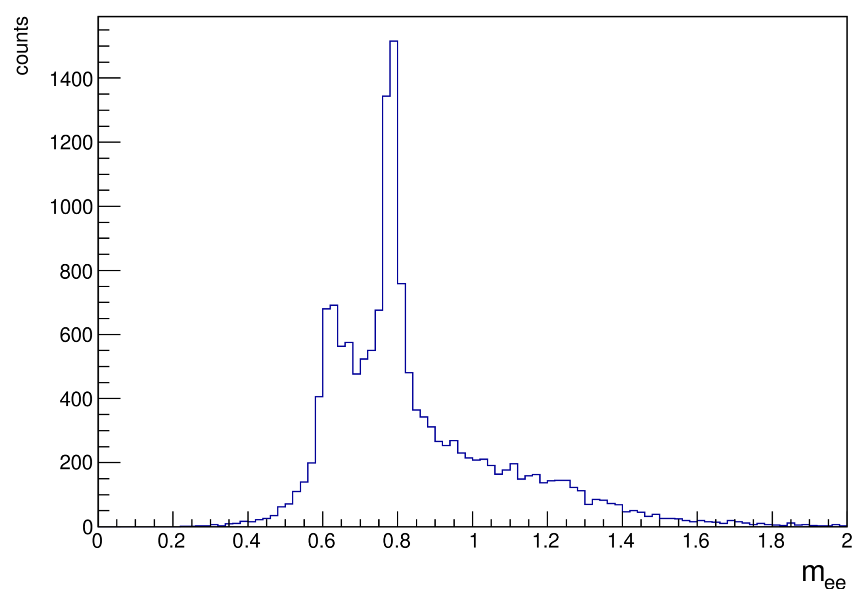
Uniform: $0.05<\mu_5<0.15$ [GeV]





Uniform: $0<\mu_5<0.2$ [GeV]





$p_{T\mu}>0.5$ GeV, $p_{T\mu\mu}>0.5$ GeV

$p_{Te}>0.15$ GeV, $p_{Tee}>0.4$ GeV

Future steps

> Full treatment of the statistical requirements and signal+background modelling

> Checking the effects of radial flow and its fluctuation (probably, other event generator will be needed)

> Full modelling of the detector response

> Analysis of real data

> Feasibility studies at NICA energy:

both theoretical (large μ_B +non-zero μ_5) and experimental/metodological

Conclusions

> There are a number of theoretical arguments and experimental evidence for the possibility of local violation of spatial parity in QCD.

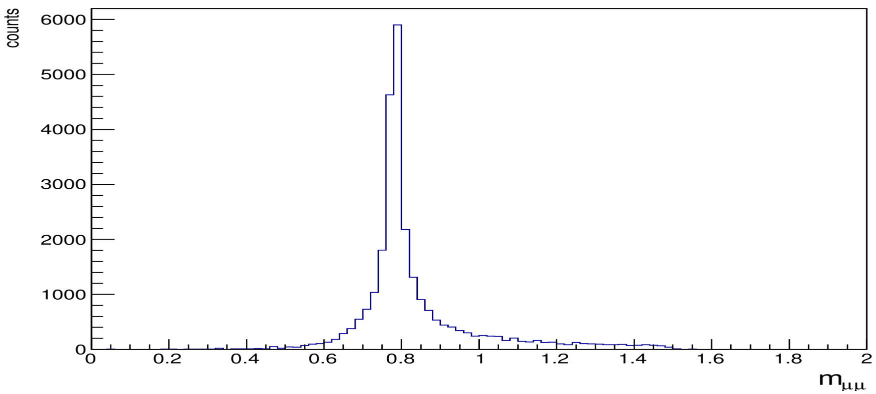
- New specific theoretical predictions need to be tested experimentally.

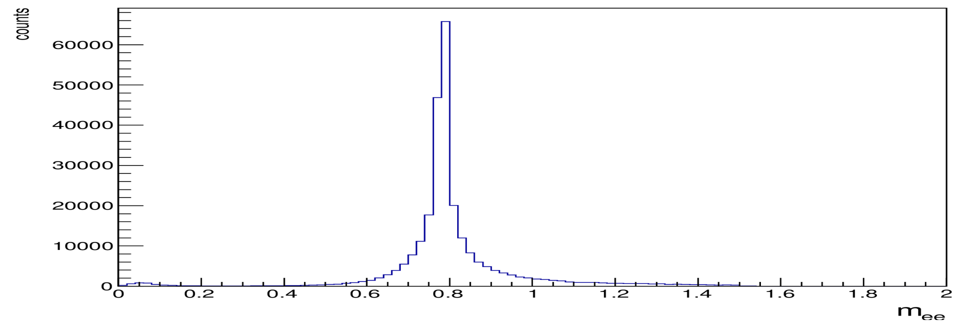
> Confirmed that angular analysis of low-mass dilepton production improve possibility of search for vector meson polarisation splitting

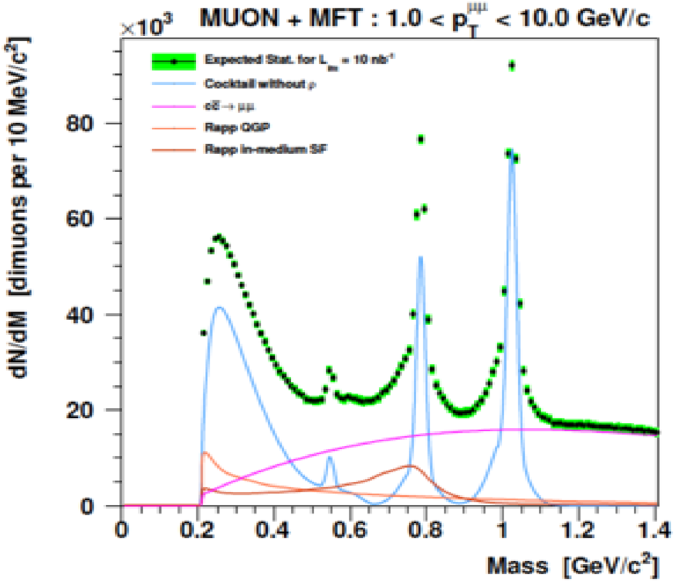
> Upgrade of the ALICE detector during the Long Shutdown 2 significantly improves the feasibility of these experimental studies at the LHC Run 3

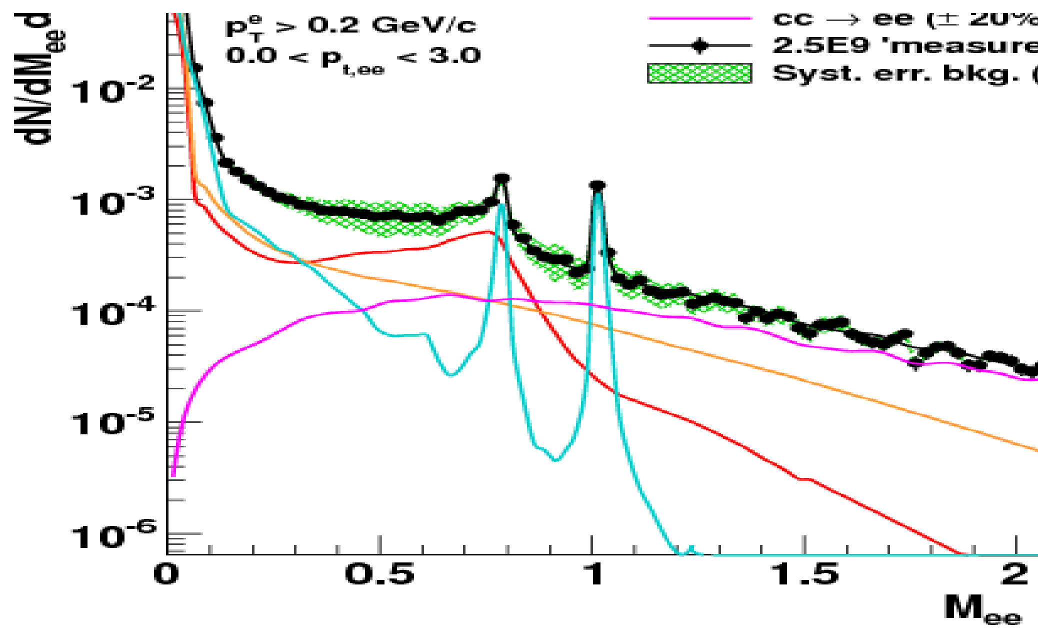
Backup

run3 conditions check mu5=0



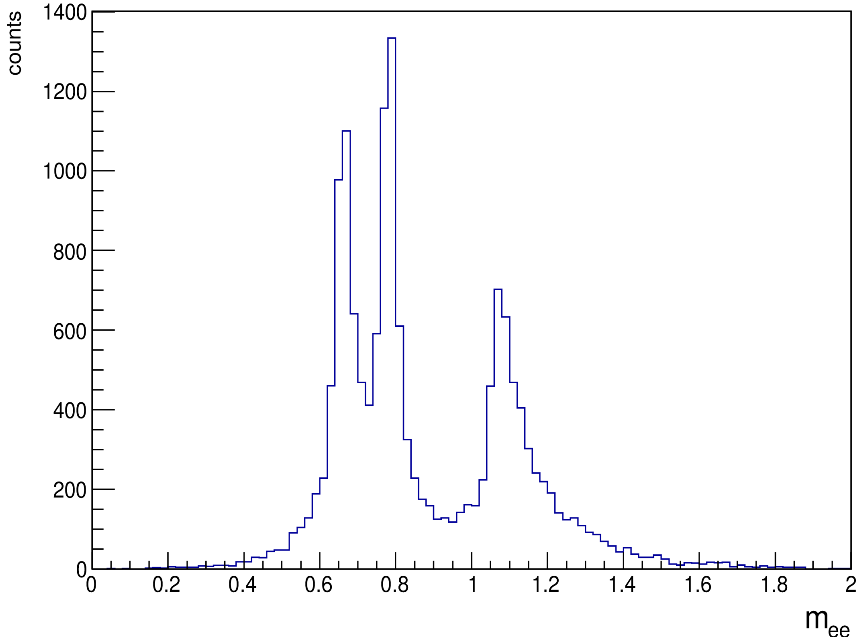


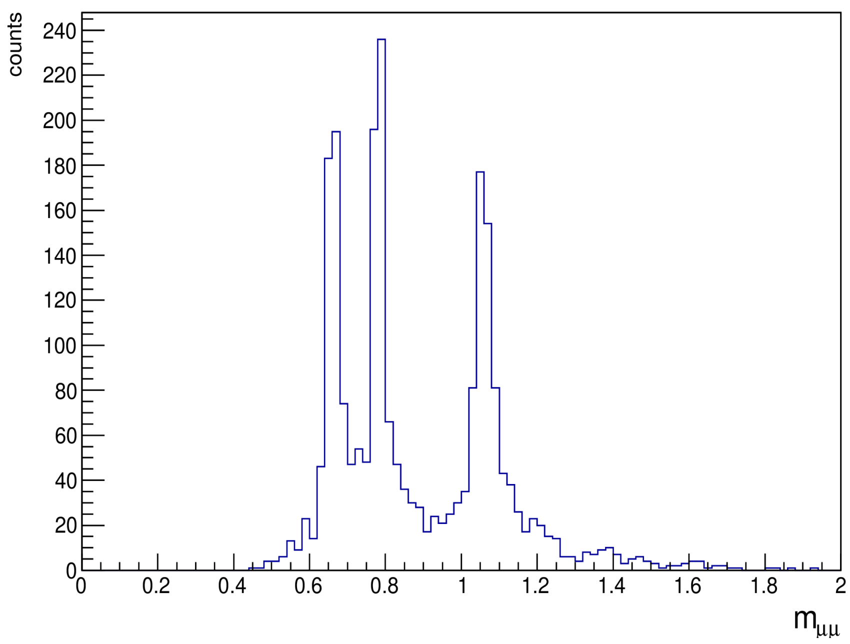




Backup

run3 conditions check mu5=0.2 GeV





CP violation in QCD

Vafa-Witten theorem: vector-like global symmetries such as parity, charge conjugation, isospin and baryon number in vector-like gauge theories like QCD cannot be spontaneously broken while the θ angle is zero

However this theorem does not apply to dense QCD matter where the partition function is not any more positive definite due to the presence of a highly non-trivial fermion determinant. In addition, out-of-equilibrium symmetry-breaking effects driven by finite temperatures are not forbidden by the Vafa-Witten theorem.

Lorentz–non-invariant P -odd operators are allowed to have non-zero expectation values at finite density $\mu > 0$ and finite temperature if the system is out of Equilibrium.

P − and CP − odd bubbles may appear in a finite volume due to large topological fluctuations in a hot medium

$$\mathcal{L}_{\mathrm{QCD}} = -\frac{1}{4} G^{\mu\nu,a} G^a_{\mu\nu} + \bar{q}(i\gamma^\mu D_\mu - m)q,$$

$$D_\mu = \partial_\mu - igG^a_\mu \lambda^a, \quad G^a_{\mu\nu} = \partial_\mu G^a_\nu - \partial_\nu G^a_\mu + g f^{abc} G^b_\mu G^c_\nu$$

$$\theta\text{-term} \qquad \Delta\mathcal{L}_\theta = \theta \frac{g^2}{16\pi^2} \mathrm{Tr}\left(G^{\mu\nu}\widetilde{G}_{\mu\nu}\right)$$

$$\theta \lesssim 10^{-9}.$$

Topological fluctuations as a source for parity breaking. Quasi-equilibrium treatment

The local partial conservation of the axial current (PCAC) relation is afflicted with the gluon anomaly

$$\partial^\mu J_{5,\mu} - 2i\bar{q}\hat{m}_q\gamma_5 q = \frac{N_f g^2}{8\pi^2}\mathrm{Tr}\left(G^{\mu\nu}\widetilde{G}_{\mu\nu}\right);$$

$$K_\mu = \frac{g^2}{2}\epsilon_{\mu\nu\rho\sigma}\mathrm{Tr}\left(G^\nu\partial^\rho G^\sigma - i\frac{2}{3}gG^\nu G^\rho G^\sigma\right), \qquad \partial_\mu K^\mu = \frac{g^2}{4}\mathrm{Tr}\left(G^{\mu\nu}\widetilde{G}_{\mu\nu}\right)$$

$$T_5 = \frac{g^2}{16\pi^2}\int_{t_i}^{t_f} dt \int_{\mathrm{vol.}} d^3x\, \mathrm{Tr}\left(G^{\mu\nu}\widetilde{G}_{\mu\nu}\right) \in \mathbb{Z},$$

$$\frac{d}{dt}(Q_5^q - 2N_f T_5) \simeq 2i\int_{\mathrm{vol.}} d^3x\,\bar{q}\hat{m}_q\gamma_5 q\ , \quad Q_5^q = \int_{\mathrm{vol.}} d^3x\,\bar{q}\gamma_0\gamma_5 q.$$

$$\langle T_5\rangle = \frac{1}{2N_f}\langle Q_5^q\rangle \quad \Longleftrightarrow \quad \mu_5 = \frac{1}{2N_f}\mu_\theta.$$