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## YIELDS AND ENERGY DISTRIBUTIONS OF α-PARTICLES IN SPONTANEOUS TERNARY NUCLEAR FISSION

Tuesday, 12 July 2022 12:30 (20 minutes)

In [1,2], a virtual mechanism of ternary fission of the nucleus (A, Z) was proposed, which is considered as a two-stage process, when at the first stage an  $\alpha$ -particle with kinetic energy  $T_{\alpha}$  close to the height of its Coulomb barrier emits from the specified nucleus, with the formation of a virtual state of the intermediate nucleus (A – 4, Z – 2), which at the second stage is involved in binary fission. Part of the energy of the emitted long-ranged  $\alpha$ -particle is taken by reducing the heat of fission of the intermediate nucleus (A – 4, Z – 2) by ( $T_{\alpha} - Q_{\alpha}$ ), where  $Q_{\alpha}$  is the heat of the true  $\alpha$ -decay of the nucleus (A, Z). The energy distribution  $W_{\alpha f}$  and yield  $N_{\alpha}$  of the  $\alpha$ -particles, taking into account the proximity of the fission widths of the nuclei (A, Z) and (A – 4, Z – 2) from the configuration (0) of these nuclei with a neck of radius  $R_{neck}$  between two fission prefragments, are defined as

W<sub>af</sub> =  $\frac{1}{2\pi} \frac{(\Gamma_{\alpha}^{A})^{0}}{(Q_{\alpha}^{A} - T_{\alpha})^{2}} = \omega_{\alpha} \frac{\hbar c \sqrt{2T_{\alpha}}}{2R_{neck}\sqrt{\mu c^{2}}} P(T_{\alpha}); N_{\alpha} = \int W_{\alpha f}(T_{\alpha}) dT_{\alpha}, where(\Gamma_{\alpha}^{A})^{0}$  is the width of the decay of the parent nucleus from the configuration (0),  $P(T_{\alpha})$  is the permeability factor of the Coulomb barrier formed by the sum of the nuclear  $V_{n}(r, \theta)$  and Coulomb  $V_{C}(r, \theta)$  interaction potentials of  $\alpha$ -particle and deformed fission fragments,  $\omega_{\alpha}$  is the probability of  $\alpha$ -particle formation in the parent nucleus,  $\mu$  is the reduced mass of ternary fission products. Calculating the permeability factor  $P(T_{\alpha})$  of the Coulomb barrier by an  $\alpha$ -particle as  $P(T_{\alpha}) = exp(-\frac{2}{\hbar c}\int_{R_{A}}^{R}\sqrt{2\mu c^{2}(V_{n}(r,\theta) + V_{C}(r,\theta)) - T_{\alpha}})dr$ , when using the deformed Coulomb potential, Saxon-Woods potential and proximity potential [4] for the nuclear potential, the energy distributions and yields of  $\alpha$ -particles for these nuclei [5 – 6].

- 1. S.G. Kadmensky et al. PEPAN 63, 620 (2022)
- 2. S.G. Kadmensky, L.V. Titova, D.E. Lyubashevsky Phys. At. Nucl. 83, 326 (2020)
- 3. L.V. Titova, Bulletin MSU. Ser. 3: Physics. Astronomy. № 5, 64 (2021)
- 4. J. Blocki, J. Randrup, W.J. Swiatecki, C.F. Tsang, Ann. Phys. (N.Y.) 105, 427 (1977)
- 5. S.Vermote et al., Nucl. Phys. A806, 1 (2008)
- 6. O.Serot, N.Carjan, C.Wagemans, Eur. Phys. J. A. 8, 187 (2000)

## The speaker is a student or young scientist

No

## Section

1. Experimental and theoretical studies of nuclear reactions

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